

# Gravitational waves in multimetric gravity

## An analysis of linearized multimetric gravity theories

Manuel Hohmann

II. Institut für theoretische Physik



Universität Hamburg  
DER FORSCHUNG | DER LEHRE | DER BILDUNG

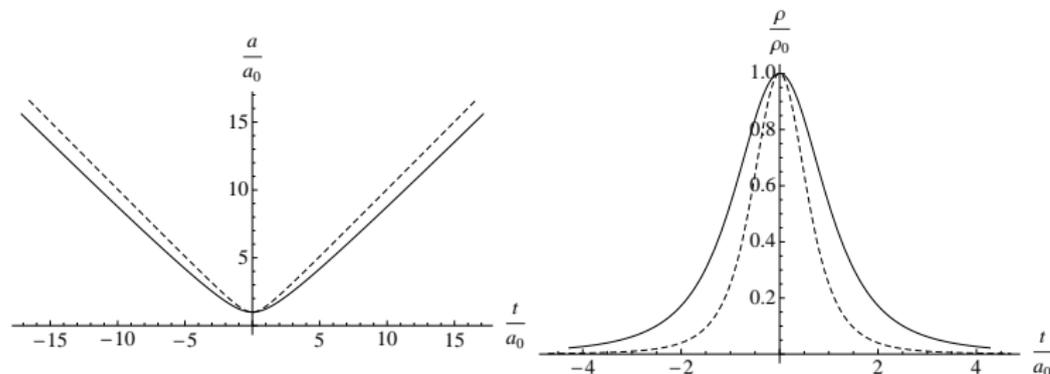
DPG-Tagung  
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# Motivation

- $\Lambda$ CDM model: 95% of the universe are dark matter / dark energy.
- Constituents of dark universe are unknown.
- Idea: DM / DE effects from additional *dark* standard model copies.
- Only interaction between standard model copies: repulsive gravity.
- Universe contains equal amounts of matter from all copies.

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- Only interaction between standard model copies: repulsive gravity.
- Universe contains equal amounts of matter from all copies.
- Dark galaxies “push” visible matter & light towards visible galaxies.  
⇒ Explanation of dark matter!
- Mutual repulsion between galaxies drives accelerating expansion.  
⇒ Explanation of dark energy! [MH, M. Wohlfarth '10]



# Construction principles

- $N \geq 2$  standard model copies  $\Psi^I$  governed by metrics  $g^I$ .
- Each standard model copy  $\Psi^I$  couples only to its own metric  $g^I$ :

$$\Rightarrow S_M[g^I, \Psi^I] = \int d^4x \sqrt{g^I} \mathcal{L}_M[g^I, \Psi^I].$$

- Different sectors couple only gravitationally:

$$\Rightarrow S = S_G[g^1, \dots, g^N] + \sum_{I=1}^N S_M[g^I, \Psi^I].$$

- Field equations obtained from variation with respect to  $g^I$ :

$$K_{ab}^I = 8\pi G_N T_{ab}^I$$

- Curvature tensor  $K_{ab}^I$  of second derivative order.
- Permutation symmetry of the sectors  $(g^I, \Psi^I)$ .
- Vacuum solution given by flat metrics  $g^I = \eta$ .

# Linearized multimetric gravity

- Perturbation ansatz:  $g^l = \eta + h^l$ .
- Most general linearized geometry tensor:

$$\underline{K}_{ab} = \underline{P} \cdot \partial^p \partial_{(a} \underline{h}_{b)p} + \underline{Q} \cdot \square \underline{h}_{ab} + \underline{R} \cdot \partial_a \partial_b \underline{h} + \underline{M} \cdot \partial^p \partial^q \underline{h}_{pq} \eta_{ab} + \underline{N} \cdot \square \underline{h} \eta_{ab}$$

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- Parameter matrices  $\underline{P}$ ,  $\underline{Q}$ ,  $\underline{R}$ ,  $\underline{M}$ ,  $\underline{N}$ .
- Permutation symmetry of the field equations:

$$P^{IJ} = (P^+ - P^-) \delta^{IJ} + P^- , \dots$$

- Simultaneously diagonalize parameter matrices.

⇒ Field equations decouple:

$$\mathfrak{K}_{ab}^1 = P_1 \partial^\rho \partial_{(a} h_{b)\rho}^1 + Q_1 \square h_{ab}^1 + R_1 \partial_a \partial_b h^1 + M_1 \partial^\rho \partial^q h_{pq}^1 \eta_{ab} + N_1 \square h^1 \eta_{ab}$$

$$\mathfrak{K}_{ab}^i = P_0 \partial^\rho \partial_{(a} h_{b)\rho}^i + Q_0 \square h_{ab}^i + R_0 \partial_a \partial_b h^i + M_0 \partial^\rho \partial^q h_{pq}^i \eta_{ab} + N_0 \square h^i \eta_{ab}$$

⇒ 10 parameters  $P_1, P_0, Q_1, Q_0, R_1, R_0, M_1, M_0, N_1, N_0$ .

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- Gravitational action:
  - Bianchi identity for  $\mathfrak{h}^1$ .
  - Geometric identity satisfied by *all* perturbations  $\mathfrak{h}^1$ . $\Rightarrow$  Restriction on parameter matrices:

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- Matter action:
  - Energy-momentum conservation:  $\nabla_a^I T^{Iab} = 0$ .
  - Field equations imply Bianchi identities for  $\mathfrak{h}^I$ .
  - Satisfied by all *solutions*  $\mathfrak{h}^I$  of the field equations.
  - $\Rightarrow$  Not a geometric identity, no restriction on parameter matrices.

# Propagation velocity

- Consider only physical degrees of freedom.
- Algebraic 3 + 1-split and differential decomposition of  $h_{ab}^I$ :
  - $4N$  scalars  $\mathfrak{I}_1^I, \mathfrak{I}_2^I, \mathfrak{I}_3^I, \mathfrak{I}_4^I$
  - $2N$  divergence-free vectors  $\mathfrak{I}_\alpha^I, \mathfrak{I}'_\alpha$
  - $N$  trace-free, divergence-free tensors  $\mathfrak{I}_{\alpha\beta}^I$
- Gauge invariance:
  - $\mathfrak{I}_3^1, \mathfrak{I}_4^1, \mathfrak{I}'_\alpha{}^1$ : pure gauge quantities
  - Remaining potentials  $\mathfrak{I}$ : gauge invariants
  - Curvature tensor  $\mathfrak{R}_{ab}^I$  depends on gauge invariants only!

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- Consider wave-like solution:

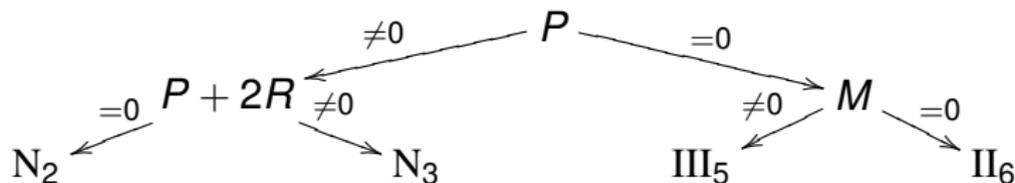
$$\mathfrak{I}(x) = \hat{\mathfrak{I}} e^{ik_a x^a}$$

⇒ Gravitational vacuum field equations solved only for  $k_a k^a = 0$ .

⇒ **Gravitational waves propagate at the speed of light!**

# Polarizations and E(2) class

- Express Riemann tensor &  $\eta^i$  in Newman-Penrose frame.
- Polarizations classified by reps. of E(2). [Eardley, Lee, Lightman *et al.* '73]
- E(2) class depends on  $P_1, R_1, M_1$  for  $\eta^1$  or  $P_0, R_0, M_0$  for  $\eta^i$ .

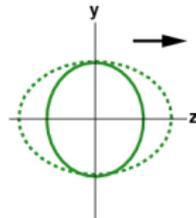
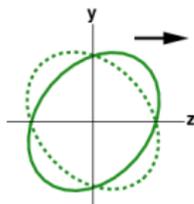
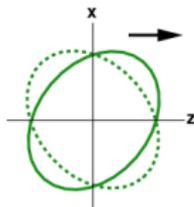
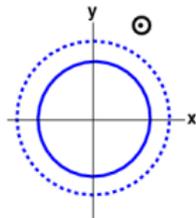
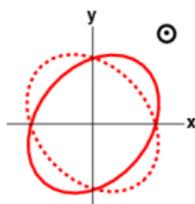
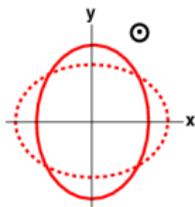


2 tensors

+1 scalar

+2 vectors

+1 scalar



# A multimetric example

Gravitational action with parameters  $x, y, u, v, w, r, s$  [MH  $\rightsquigarrow$  PRD 03/12]:

$$S_G = \frac{1}{16\pi} \int d^4x \sqrt{g_0} \left[ x \sum_{I,J=1}^N g^{Iij} R^J_{ij} + \sum_{I=1}^N g^{Iij} \left( y R^I_{ij} + u \tilde{S}^I_i \tilde{S}^I_j \right. \right. \\ \left. \left. + v \tilde{S}^I_k \tilde{S}^Ik_{ij} + w \tilde{S}^Ik_{im} \tilde{S}^Im_{jk} + g^{Ikl} g^{lmn} \left( r \tilde{S}^Im_{ik} \tilde{S}^In_{jl} + s \tilde{S}^Im_{ij} \tilde{S}^In_{kl} \right) \right) \right].$$

Connection difference tensors:

$$S^{Ji}_{jk} = \Gamma^{Ii}_{jk} - \Gamma^{Ji}_{jk}, \quad S^J_j = S^{Jk}_{jk}, \\ \tilde{S}^{Ji}_{jk} = \frac{1}{N} \sum_{I=1}^N S^{Ii}_{jk}, \quad \tilde{S}^J_j = \tilde{S}^{Jk}_{jk}.$$

Volume form:

$$g_0 = \prod_{I=1}^N (g^I)^{\frac{1}{N}}$$

# Parameter matrices

- Calculate linearized field equations.
- Eigenvalues of the parameter matrices:

$$\begin{aligned} P_1 &= -2Q_1 = -2R_1 = -2M_1 = 2N_1 = Nx + y, \\ P_0 &= -Nx + y - w + r - 2s, \quad R_0 = M_0 = \frac{Nx - v + 2s}{2}, \\ Q_0 &= \frac{Nx - y + w - 3r}{2}, \quad N_0 = \frac{-Nx - y - u + v - s}{2}. \end{aligned}$$

⇒ Field equations are gauge invariant. ✓

⇒ Gravitational action is diffeomorphism invariant. ✓

# Parametrized post-Newtonian (PPN) formalism

- Obtain “fingerprint” of metric gravity theories. [Thorne, Will '71; Will '93]
  - Characterize gravity theories by 10 parameters.
  - PPN parameters can be measured by solar system experiments.
- Extension to multimetric gravity theories. [MH, M. Wohlfarth '10]
  - Additional 14 unobserved parameters.
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  - Additional 14 unobserved parameters.
  - 8 parameters can be obtained from linearized field equations.
- PPN formalism applied to example theory.
- Restriction of input parameters:

$$y = \frac{1}{2 - N} - Nx, \quad v = \frac{6 - N}{4 - 2N} - Nx + 2u,$$
$$w = -\frac{6 - N}{4 - 2N} + Nx - 3u, \quad r = -\frac{1}{2 - N} + Nx - u.$$

⇒ Three free parameters  $x$ ,  $u$ ,  $s$  remain.

⇒ Consistent with solar system experiments up to linear PPN order.

# PPN consistent multimetric example

Parameter matrices:

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- Dependent only on single parameter  $p := Nx - u + s$ .

Generic case:

- E(2) class for  $\mathfrak{h}^1$ :  $N_2$
- E(2) class for  $\mathfrak{h}^i$ :  $N_2$

⇒ Effective E(2) class:  $N_2$

Special case  $p = \frac{6-N}{8-4N}$ :

- E(2) class for  $\mathfrak{h}^1$ :  $N_2$
- E(2) class for  $\mathfrak{h}^i$ :  $\text{II}_6$

⇒ Effective E(2) class:  $\text{II}_6$

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- **Example theory:**
  - Defined by gravitational action
  - Consistent up to linear PPN level
  - 3 free parameters
  - $E(2)$  class either  $N_2$  or  $II_6$

- Production of gravitational waves:
  - Slow-moving sources
  - Binary systems
    - Expansion in multipole coefficients
    - Make use of tensor solid harmonics  $\Rightarrow$  algebraic equations
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$\rightsquigarrow$  **Experimental data from gravitational wave experiments!**