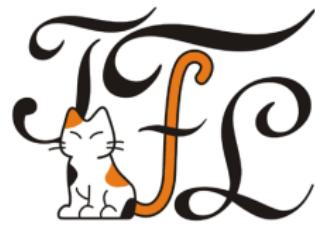


Parameterized post-Newtonian limit of Horndeski's gravity theory

Phys. Rev. D **92** 064019

Manuel Hohmann

Laboratory of Theoretical Physics - Institute of Physics - University of Tartu
Center of Excellence "The Dark Side of the Universe"



DPG Spring Conference - Session GR 5 - 1. March 2016

Motivation

- So far unexplained cosmological observations:
 - Accelerating expansion of the universe
 - Homogeneity of cosmic microwave background

Motivation

- So far unexplained cosmological observations:
 - Accelerating expansion of the universe
 - Homogeneity of cosmic microwave background
- Models for explaining these observations:
 - Λ CDM model / dark energy
 - Inflation

Motivation

- So far unexplained cosmological observations:
 - Accelerating expansion of the universe
 - Homogeneity of cosmic microwave background
- Models for explaining these observations:
 - Λ CDM model / dark energy
 - Inflation
- Physical mechanisms are not understood:
 - Unknown type of matter?
 - Modification of the laws of gravity?
 - Scalar field in addition to metric mediating gravity?
 - Quantum gravity effects?

Motivation

- So far unexplained cosmological observations:
 - Accelerating expansion of the universe
 - Homogeneity of cosmic microwave background
- Models for explaining these observations:
 - Λ CDM model / dark energy
 - Inflation
- Physical mechanisms are not understood:
 - Unknown type of matter?
 - Modification of the laws of gravity?
 - **Scalar field in addition to metric mediating gravity?**
 - Quantum gravity effects?
- **Horndeski gravity** [G. W. Horndeski '74]:
 - Scalar-tensor theory of gravity.
 - Most general STG with second order field equations.
 - Healthy, ghost-free theory.
 - Contains many interesting cases (Galileons, Higgs inflation...).

Gravitational action

- Action functional [T. Kobayashi, M. Yamaguchi, J. 'i. Yokoyama '11]:

$$S = \sum_{i=2}^5 \int d^4x \sqrt{-g} \mathcal{L}_i[g_{\mu\nu}, \phi] + S_m[g_{\mu\nu}, \chi_m].$$

Gravitational action

- Action functional [T. Kobayashi, M. Yamaguchi, J. 'i. Yokoyama '11]:

$$S = \sum_{i=2}^5 \int d^4x \sqrt{-g} \mathcal{L}_i[g_{\mu\nu}, \phi] + S_m[g_{\mu\nu}, \chi_m].$$

- Gravitational Lagrangian with $X = -\nabla_\mu \phi \nabla^\mu \phi$:

$$\mathcal{L}_2 = K(\phi, X),$$

$$\mathcal{L}_3 = -G_3(\phi, X) \square \phi,$$

$$\mathcal{L}_4 = G_4(\phi, X) R + G_{4X}(\phi, X) \left[(\square \phi)^2 - (\nabla_\mu \nabla_\nu \phi)^2 \right],$$

$$\mathcal{L}_5 = G_5(\phi, X) G_{\mu\nu} \nabla^\mu \nabla^\nu \phi$$

$$-\frac{1}{6} G_{5X}(\phi, X) \left[(\square \phi)^3 - 3(\square \phi)(\nabla_\mu \nabla_\nu \phi)^2 + 2(\nabla_\mu \nabla_\nu \phi)^3 \right].$$

Gravitational action

- Action functional [T. Kobayashi, M. Yamaguchi, J. 'i. Yokoyama '11]:

$$S = \sum_{i=2}^5 \int d^4x \sqrt{-g} \mathcal{L}_i[g_{\mu\nu}, \phi] + S_m[g_{\mu\nu}, \chi_m].$$

- Gravitational Lagrangian with $X = -\nabla_\mu \phi \nabla^\mu \phi$:

$$\mathcal{L}_2 = K(\phi, X),$$

$$\mathcal{L}_3 = -G_3(\phi, X) \square \phi,$$

$$\mathcal{L}_4 = G_4(\phi, X) R + G_{4X}(\phi, X) \left[(\square \phi)^2 - (\nabla_\mu \nabla_\nu \phi)^2 \right],$$

$$\mathcal{L}_5 = G_5(\phi, X) G_{\mu\nu} \nabla^\mu \nabla^\nu \phi$$

$$-\frac{1}{6} G_{5X}(\phi, X) \left[(\square \phi)^3 - 3(\square \phi)(\nabla_\mu \nabla_\nu \phi)^2 + 2(\nabla_\mu \nabla_\nu \phi)^3 \right].$$

- Free functions K, G_3, G_4, G_5 .

Perturbative expansion

- Background solution:
 - Minkowski metric $\eta_{\mu\nu}$
 - Constant scalar field value Φ

Perturbative expansion

- Background solution:
 - Minkowski metric $\eta_{\mu\nu}$
 - Constant scalar field value Φ
- Perturbation of dynamical fields:

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad \phi = \Phi + \psi, \quad X = -\frac{1}{2}\nabla^\mu\psi\nabla_\mu\psi.$$

Perturbative expansion

- Background solution:
 - Minkowski metric $\eta_{\mu\nu}$
 - Constant scalar field value Φ
- Perturbation of dynamical fields:

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad \phi = \Phi + \psi, \quad X = -\frac{1}{2}\nabla^\mu\psi\nabla_\mu\psi.$$

- Taylor expansion of free functions:

$$K(\phi, X) = \sum_{m,n=0}^{\infty} K_{(m,n)} \psi^m X^n.$$

- Expansion coefficients:

$$K_{(m,n)} = \frac{1}{m!n!} \left. \frac{\partial^{m+n}}{\partial\phi^m\partial X^n} K(\phi, X) \right|_{\phi=\Phi, X=0}.$$

- Similar expansion for G_3, G_4, G_5 .

Post-Newtonian approximation

- Perfect fluid energy-momentum tensor:

$$T^{\mu\nu} = (\rho + p\Pi)u^\mu u^\nu + pg^{\mu\nu}.$$

- Four-velocity u^μ .
- Matter density ρ .
- Specific internal energy Π .
- Pressure p .

Post-Newtonian approximation

- Perfect fluid energy-momentum tensor:

$$T^{\mu\nu} = (\rho + p\Pi) u^\mu u^\nu + p g^{\mu\nu}.$$

- Four-velocity u^μ .
- Matter density $\rho \sim \mathcal{O}(2)$.
- Specific internal energy $\Pi \sim \mathcal{O}(2)$.
- Pressure $p \sim \mathcal{O}(4)$.
- Slow-moving source matter:
$$v^i = \frac{u^i}{u^0} \ll 1.$$
- Assign velocity orders $|v^i|^n \sim \mathcal{O}(n)$ based on solar system.

Post-Newtonian approximation

- Perfect fluid energy-momentum tensor:

$$T^{\mu\nu} = (\rho + p\Pi)u^\mu u^\nu + pg^{\mu\nu}.$$

- Four-velocity u^μ .
- Matter density $\rho \sim \mathcal{O}(2)$.
- Specific internal energy $\Pi \sim \mathcal{O}(2)$.
- Pressure $p \sim \mathcal{O}(4)$.
- Slow-moving source matter:

$$v^i = \frac{u^i}{u^0} \ll 1.$$

- Assign velocity orders $|v^i|^n \sim \mathcal{O}(n)$ based on solar system.
- Relevant terms for dynamical fields:

$$h_{00}^{(2)}, \quad h_{ij}^{(2)}, \quad h_{0j}^{(3)}, \quad h_{00}^{(4)}, \quad \psi^{(2)}, \quad \psi^{(4)}.$$

Post-Newtonian approximation

- Perfect fluid energy-momentum tensor:

$$T^{\mu\nu} = (\rho + p\Pi)u^\mu u^\nu + pg^{\mu\nu}.$$

- Four-velocity u^μ .
- Matter density $\rho \sim \mathcal{O}(2)$.
- Specific internal energy $\Pi \sim \mathcal{O}(2)$.
- Pressure $p \sim \mathcal{O}(4)$.
- Slow-moving source matter:

$$v^i = \frac{u^i}{u^0} \ll 1.$$

- Assign velocity orders $|v^i|^n \sim \mathcal{O}(n)$ based on solar system.
- Relevant terms for dynamical fields:

$$h_{00}^{(2)}, \quad h_{ij}^{(2)}, \quad h_{0j}^{(3)}, \quad h_{00}^{(4)}, \quad \psi^{(2)}, \quad \psi^{(4)}.$$

- Time dependence only through motion of source matter.
⇒ Assign time derivative $\partial_0 \sim \mathcal{O}(1)$.

Spherically symmetric solution

- Static, point-like mass source:

$$\rho = M\delta(\vec{x}), \quad \Pi = 0, \quad p = 0, \quad v_i = 0.$$

Spherically symmetric solution

- Static, point-like mass source:

$$\rho = M\delta(\vec{x}), \quad \Pi = 0, \quad p = 0, \quad v_i = 0.$$

- Spherically symmetric metric:

$$g_{00} = -1 + 2G_{\text{eff}}(r)U(r) - 2G_{\text{eff}}^2(r)\beta(r)U^2(r) + \Phi^{(4)}(r) + \mathcal{O}(6),$$

$$g_{0j} = \mathcal{O}(5),$$

$$g_{ij} = [1 + 2G_{\text{eff}}(r)\gamma(r)U(r)]\delta_{ij} + \mathcal{O}(4).$$

- Newtonian potential: $U(r) = M/r$.
- Gravitational self energy $\Phi^{(4)}(r)$.
- Effective gravitational constant $G_{\text{eff}}(r)$.
- PPN parameters $\gamma(r)$ and $\beta(r)$.

Spherically symmetric solution

- Static, point-like mass source:

$$\rho = M\delta(\vec{x}), \quad \Pi = 0, \quad p = 0, \quad v_i = 0.$$

- Spherically symmetric metric:

$$g_{00} = -1 + 2G_{\text{eff}}(r)U(r) - 2G_{\text{eff}}^2(r)\beta(r)U^2(r) + \Phi^{(4)}(r) + \mathcal{O}(6),$$

$$g_{0j} = \mathcal{O}(5),$$

$$g_{ij} = [1 + 2G_{\text{eff}}(r)\gamma(r)U(r)]\delta_{ij} + \mathcal{O}(4).$$

- Newtonian potential: $U(r) = M/r$.
- Gravitational self energy $\Phi^{(4)}(r)$.
- Effective gravitational constant $G_{\text{eff}}(r)$.
- PPN parameters $\gamma(r)$ and $\beta(r)$.

- Consistency condition:

$$K_{(0,0)} = K_{(1,0)} = 0.$$

Scalar field $\psi^{(2)}$

- Scalar field equation at $\mathcal{O}(2)$ is screened Poisson equation:

$$\psi_{,ii}^{(2)} - m_\psi^2 \psi^{(2)} = -c_\psi \rho.$$

Scalar field $\psi^{(2)}$

- Scalar field equation at $\mathcal{O}(2)$ is screened Poisson equation:

$$\psi_{,ii}^{(2)} - m_\psi^2 \psi^{(2)} = -c_\psi \rho.$$

- Solution:

$$\psi^{(2)}(r) = \frac{M}{4\pi r} c_\psi e^{-m_\psi r}.$$

Scalar field $\psi^{(2)}$

- Scalar field equation at $\mathcal{O}(2)$ is screened Poisson equation:

$$\psi_{,ii}^{(2)} - \textcolor{red}{m}_\psi^2 \psi^{(2)} = -\textcolor{red}{c}_\psi \rho.$$

- Solution:

$$\psi^{(2)}(r) = \frac{M}{4\pi r} \textcolor{red}{c}_\psi e^{-\textcolor{red}{m}_\psi r}.$$

- Constants:

$$\textcolor{red}{m}_\psi = \sqrt{\frac{-2K_{(2,0)}}{K_{(0,1)} - 2G_{3(1,0)} + 3\frac{G_{4(1,0)}^2}{G_{4(0,0)}}}},$$

$$\textcolor{red}{c}_\psi = \frac{G_{4(1,0)}}{2G_{4(0,0)}} \left(K_{(0,1)} - 2G_{3(1,0)} + 3\frac{G_{4(1,0)}^2}{G_{4(0,0)}} \right)^{-1}.$$

Effective gravitational constant $G_{\text{eff}}(r)$

- Metric field equation:

$$h_{00,ii}^{(2)} = c_1 \psi^{(2)} - c_2 \rho .$$

Effective gravitational constant $G_{\text{eff}}(r)$

- Metric field equation:

$$h_{00,ii}^{(2)} = c_1 \psi^{(2)} - c_2 \rho.$$

- Solve and read off effective gravitational constant:

$$G_{\text{eff}}(r) = \frac{1}{8\pi} \left[c_2 + \frac{c_1 c_\psi}{m_\psi^2} (e^{-m_\psi r} - 1) \right].$$

Effective gravitational constant $G_{\text{eff}}(r)$

- Metric field equation:

$$h_{00,ii}^{(2)} = \textcolor{red}{c}_1 \psi^{(2)} - \textcolor{red}{c}_2 \rho.$$

- Solve and read off effective gravitational constant:

$$G_{\text{eff}}(r) = \frac{1}{8\pi} \left[\textcolor{red}{c}_2 + \frac{\textcolor{red}{c}_1 c_\psi}{m_\psi^2} (e^{-m_\psi r} - 1) \right].$$

- Constants:

$$\textcolor{red}{c}_1 = -2 \frac{G_{4(1,0)} K_{(2,0)}}{G_{4(0,0)}} \left(K_{(0,1)} - 2G_{3(1,0)} + 3 \frac{G_{4(1,0)}^2}{G_{4(0,0)}} \right)^{-1},$$

$$\textcolor{red}{c}_2 = \frac{1}{G_{4(0,0)}} \left[\frac{1}{2} + \frac{G_{4(1,0)}^2}{2G_{4(0,0)}} \left(K_{(0,1)} - 2G_{3(1,0)} + 3 \frac{G_{4(1,0)}^2}{G_{4(0,0)}} \right)^{-1} \right].$$

PPN parameter γ

- Metric field equation:

$$h_{ij,kk}^{(2)} = \left(c_3 \psi^{(2)} - c_4 \rho \right) \delta_{ij} .$$

PPN parameter γ

- Metric field equation:

$$h_{ij,kk}^{(2)} = \left(c_3 \psi^{(2)} - c_4 \rho \right) \delta_{ij} .$$

- Solve and read off PPN parameter γ :

$$\gamma(r) = \frac{c_4 + \frac{c_3 c_\psi}{m_\psi^2} (e^{-m_\psi r} - 1)}{c_2 + \frac{c_1 c_\psi}{m_\psi^2} (e^{-m_\psi r} - 1)} .$$

PPN parameter γ

- Metric field equation:

$$h_{ij,kk}^{(2)} = \left(\textcolor{red}{c}_3 \psi^{(2)} - \textcolor{red}{c}_4 \rho \right) \delta_{ij}.$$

- Solve and read off PPN parameter γ :

$$\gamma(r) = \frac{\textcolor{red}{c}_4 + \frac{c_3 c_\psi}{m_\psi^2} (e^{-m_\psi r} - 1)}{c_2 + \frac{c_1 c_\psi}{m_\psi^2} (e^{-m_\psi r} - 1)}.$$

- Constants:

$$\textcolor{red}{c}_3 = 2 \frac{G_{4(1,0)} K_{(2,0)}}{G_{4(0,0)}} \left(K_{(0,1)} - 2G_{3(1,0)} + 3 \frac{G_{4(1,0)}^2}{G_{4(0,0)}} \right)^{-1},$$

$$\textcolor{red}{c}_4 = \frac{1}{G_{4(0,0)}} \left[\frac{1}{2} - \frac{G_{4(1,0)}^2}{2G_{4(0,0)}} \left(K_{(0,1)} - 2G_{3(1,0)} + 3 \frac{G_{4(1,0)}^2}{G_{4(0,0)}} \right)^{-1} \right].$$

PPN parameter γ

- Metric field equation:

$$h_{ij,kk}^{(2)} = \left(c_3 \psi^{(2)} - c_4 \rho \right) \delta_{ij} .$$

- Solve and read off PPN parameter γ :

$$\gamma(r) = \frac{2\omega + 3 - e^{-m_\psi r}}{2\omega + 3 + e^{-m_\psi r}} .$$

- Constants:

$$\omega = \frac{G_{4(0,0)}}{2G_{4(1,0)}^2} (K_{(0,1)} - 2G_{3(1,0)}) ,$$

$$m_\psi = \sqrt{\frac{-2K_{(2,0)}}{K_{(0,1)} - 2G_{3(1,0)} + 3\frac{G_{4(1,0)}^2}{G_{4(0,0)}}}} .$$

PPN parameter β

- Calculate β from fourth order solution:

$$\begin{aligned}\beta(r) = 1 + \frac{1}{(2\omega + 3 + e^{-m_\psi r})^2} & \left\{ \frac{\omega + \tau - 4\omega\sigma}{2\omega + 3} e^{-2m_\psi r} \right. \\ & + (2\omega + 3)m_\psi r \left[e^{-m_\psi r} \ln(m_\psi r) - \frac{1}{2}e^{-2m_\psi r} \right. \\ & \quad \left. \left. - (m_\psi r + e^{m_\psi r}) \operatorname{Ei}(-2m_\psi r) \right] \right. \\ & + \frac{6\mu r + 3(3\omega + \tau + 6\sigma + 3)m_\psi^2 r}{2(2\omega + 3)m_\psi} \left[e^{m_\psi r} \operatorname{Ei}(-3m_\psi r) \right. \\ & \quad \left. \left. - e^{-m_\psi r} \operatorname{Ei}(-m_\psi r) \right] \right\},\end{aligned}$$

PPN parameter β

- Calculate β from fourth order solution:

$$\begin{aligned}\beta(r) = & 1 + \frac{1}{(2\omega + 3 + e^{-m_\psi r})^2} \left\{ \frac{\omega + \tau - 4\omega\sigma}{2\omega + 3} e^{-2m_\psi r} \right. \\ & + (2\omega + 3)m_\psi r \left[e^{-m_\psi r} \ln(m_\psi r) - \frac{1}{2}e^{-2m_\psi r} \right. \\ & \quad \left. \left. - (m_\psi r + e^{m_\psi r}) \operatorname{Ei}(-2m_\psi r) \right] \right. \\ & + \frac{6\mu r + 3(3\omega + \tau + 6\sigma + 3)m_\psi^2 r}{2(2\omega + 3)m_\psi} [e^{m_\psi r} \operatorname{Ei}(-3m_\psi r) \\ & \quad \left. \left. - e^{-m_\psi r} \operatorname{Ei}(-m_\psi r) \right] \right\},\end{aligned}$$

- Constants $m_\psi, \omega, \tau, \sigma, \mu$.

Limiting cases

- $m_\psi \rightarrow 0$, all other constants fixed and finite:

$$\gamma = \frac{\omega + 1}{\omega + 2}, \quad \beta = 1 + \frac{\omega + \tau - 4\omega\sigma}{(2\omega + 3)(2\omega + 4)^2}.$$

Limiting cases

- $m_\psi \rightarrow 0$, all other constants fixed and finite:

$$\gamma = \frac{\omega + 1}{\omega + 2}, \quad \beta = 1 + \frac{\omega + \tau - 4\omega\sigma}{(2\omega + 3)(2\omega + 4)^2}.$$

- $\omega \rightarrow \infty$, all other constants fixed and finite:

$$\gamma = \beta = 1.$$

Limiting cases

- $m_\psi \rightarrow 0$, all other constants fixed and finite:

$$\gamma = \frac{\omega + 1}{\omega + 2}, \quad \beta = 1 + \frac{\omega + \tau - 4\omega\sigma}{(2\omega + 3)(2\omega + 4)^2}.$$

- $\omega \rightarrow \infty$, all other constants fixed and finite:

$$\gamma = \beta = 1.$$

- $m_\psi r \rightarrow \infty$, large distance from the matter source:

$$\gamma = \beta = 1.$$

Example: scalar-tensor gravity with potential

- Gravitational action:

$$S_G = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} \left(\phi R - \frac{\omega(\phi)}{\phi} \partial_\rho \phi \partial^\rho \phi - 2\kappa^2 V(\phi) \right).$$

Example: scalar-tensor gravity with potential

- Gravitational action:

$$S_G = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} \left(\phi R - \frac{\omega(\phi)}{\phi} \partial_\rho \phi \partial^\rho \phi - 2\kappa^2 V(\phi) \right).$$

- PPN parameters [MH, L. Järv, P. Kuusk, E. Randla '13]:

$$\gamma(r) = \frac{2\omega_0 + 3 - e^{-m_\psi r}}{2\omega_0 + 3 + e^{-m_\psi r}},$$

$$\beta(r) = 1 + \frac{1}{(2\omega_0 + 3 + e^{-m_\psi r})^2} \left\{ \frac{\Phi\omega_1}{2\omega_0 + 3} e^{-2m_\psi r} + (2\omega_0 + 3)m_\psi r \right.$$

$$\times \left[e^{-m_\psi r} \ln(m_\psi r) - (m_\psi r + e^{m_\psi r}) \operatorname{Ei}(-2m_\psi r) - \frac{1}{2} e^{-2m_\psi r} \right]$$

$$+ \frac{3m_\psi r}{2} \left(1 - \frac{\Phi V_3}{V_2} + \frac{\Phi\omega_1}{2\omega_0 + 3} \right) [e^{m_\psi r} \operatorname{Ei}(-3m_\psi r) - e^{-m_\psi r} \operatorname{Ei}(-m_\psi r)] \right\}.$$

Example: Higgs inflation

- Gravitational action [F. L. Bezrukov, M. Shaposhnikov '08]:

$$S_G = \int d^4x \sqrt{-g} \left(\frac{M_{\text{Pl}}^2 - \xi \phi^2}{2} R + X - V(\phi) \right).$$

Example: Higgs inflation

- Gravitational action [F. L. Bezrukov, M. Shaposhnikov '08]:

$$S_G = \int d^4x \sqrt{-g} \left(\frac{M_{\text{Pl}}^2 - \xi \phi^2}{2} R + X - V(\phi) \right).$$

- PPN parameters:

$$\gamma = 1 - 4\xi^2 e^{-m_\psi r} \frac{\Phi^2}{M_{\text{Pl}}^2} + \mathcal{O}\left(\frac{\Phi^3}{M_{\text{Pl}}^3}\right),$$

$$\begin{aligned} \beta = 1 &+ \left\{ 2\xi^3 e^{-2m_\psi r} - \xi^2 m_\psi r [e^{-2m_\psi r} - 2e^{-m_\psi r} \ln(m_\psi r) \right. \\ &\left. + 2(m_\psi r + e^{m_\psi r}) \text{Ei}(-2m_\psi r)] \right\} \frac{\Phi^2}{M_{\text{Pl}}^2} + \mathcal{O}\left(\frac{\Phi^3}{M_{\text{Pl}}^3}\right). \end{aligned}$$

Summary

- Horndeski's gravity theory:
 - Most general scalar-tensor theory with second order equations.
 - Four free functions of ϕ and $X = -\nabla^\mu \phi \nabla_\mu \phi / 2$.

Summary

- Horndeski's gravity theory:
 - Most general scalar-tensor theory with second order equations.
 - Four free functions of ϕ and $X = -\nabla^\mu \phi \nabla_\mu \phi / 2$.
- Example theories:
 - Classical scalar-tensor gravity with arbitrary potential.
 - Models of Higgs inflation.
 - Galileons.

Summary

- Horndeski's gravity theory:
 - Most general scalar-tensor theory with second order equations.
 - Four free functions of ϕ and $X = -\nabla^\mu \phi \nabla_\mu \phi / 2$.
- Example theories:
 - Classical scalar-tensor gravity with arbitrary potential.
 - Models of Higgs inflation.
 - Galileons.
- PPN parameters:
 - Most general theory: obtained PPN parameters $\gamma(r)$ and $\beta(r)$.
 - Massless scalar field: only γ and β potentially deviate.
 - Reproduces and generalizes well-known results.

Outlook

- Extend analysis to more general theories:
 - Allow time-dependent scalar background field $\dot{\phi} \neq 0$.
 - Theories beyond Horndeski / G^3 -inflation.
 - Multi-scalar Horndeski gravity.

Outlook

- Extend analysis to more general theories:
 - Allow time-dependent scalar background field $\dot{\phi} \neq 0$.
 - Theories beyond Horndeski / G^3 -inflation.
 - Multi-scalar Horndeski gravity.
- Take screening mechanisms into account:
 - Vainshtein mechanism.
 - Chameleon mechanism.
 - Symmetron mechanism.

Acknowledgments

- Estonian Research Council
 - ERMOS115 “Geometric extensions of general relativity - foundations and phenomenology”
 - PUT790 “Geometric foundations of gravity and their comparison to observations”
- Archimedes Foundation
 - TK133 “The Dark Side of the Universe”

Job openings supported by these grants:

- PhD student position (deadline: 20. March 2016)
- PostDoc position (deadline: 10. April 2016)

<http://www.fi.ut.ee/en/postdoc-and-phd-in-gravity-theory>