

Generalized scalar-torsion theories of gravity

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Teleparallel gravity workshop, Tartu 25. - 29. June 2018

Motivation

- Open questions in cosmology and gravity:
 - Accelerating phases in the history of the Universe?
 - Relation between gravity and gauge theories?
 - How to quantize gravity?

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 - Describes gravity as gauge theory of the translation group.
 - Gravitational field strength is torsion.
 - First order action, second order field equations.
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 - Possibly arises from more fundamental theory.
 - Differs from non-minimal coupling to curvature.
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- Scalar field non-minimally coupled to torsion [Geng '11]:
 - Possibly arises from more fundamental theory.
 - Differs from non-minimal coupling to curvature.
 - Possible model for so far unexplained observations.
- Arising questions:
 - Most general class of scalar-torsion gravity theories?
 - Behavior under conformal transformations?

Ingredients of scalar-torsion gravity

- Fundamental fields:

- Coframe field $\theta^a = \theta^a_{\mu} dx^{\mu}$.
- Flat spin connection $\overset{\bullet}{\omega}{}^a{}_b = \overset{\bullet}{\omega}{}^a{}_{b\mu} dx^{\mu}$.
- N scalar fields $\phi = (\phi^A; A = 1, \dots, N)$.
- Arbitrary matter fields χ^I .

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- Derived quantities:

- Frame field $e_a = e_a{}^\mu \partial_\mu$ with $\iota_{e_a} \theta^b = \delta_a^b$.
- Metric $g_{\mu\nu} = \eta_{ab} \theta^a{}_\mu \theta^b{}_\nu$.
- Volume form $\theta d^4x = \theta^0 \wedge \theta^1 \wedge \theta^2 \wedge \theta^3$.
- Levi-Civita connection

$$\overset{\circ}{\omega}_{ab} = -\frac{1}{2}(\iota_{e_b} \iota_{e_c} d\theta_a + \iota_{e_c} \iota_{e_a} d\theta_b - \iota_{e_a} \iota_{e_b} d\theta_c)\theta^c.$$

- Torsion $T^a = d\theta^a + \overset{\bullet}{\omega}{}^a{}_b \wedge \theta^b$.

Overview

$$S_g [\theta^a, \overset{\bullet}{\omega}{}^a{}_b, \phi^A] + S_m [\theta^a, \phi^A, \chi^I]$$

$$L(T, X, Y, \phi)$$

$$\mathcal{A}(\phi)T + \mathcal{B}(\phi)\partial_\mu\phi\partial^\mu\phi + \mathcal{C}(\phi)T^\mu\partial_\mu\phi + \mathcal{V}(\phi)$$

$$\mathcal{A}(\phi)T + \mathcal{B}(\phi)\partial_\mu\phi\partial^\mu\phi + \mathcal{V}(\phi)$$

$$f(T, \phi) + Z(\phi)\partial_\mu\phi\partial^\mu\phi$$

General scalar-torsion gravity - action

- Structure of the action [MH '18]:

$$S \left[\theta^a, \overset{\bullet}{\omega}{}^a{}_b, \phi^A, \chi^I \right] = S_g \left[\theta^a, \overset{\bullet}{\omega}{}^a{}_b, \phi^A \right] + S_m \left[\theta^a, \phi^A, \chi^I \right].$$

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- Variation of the action:
 - Gravitational part:

$$\begin{aligned}\delta S_g &= \int_M \left(\Delta_a \wedge \delta \theta^a + \frac{1}{2} \Xi_a{}^b \wedge \delta \overset{\bullet}{\omega}{}^a{}_b + \Phi_A \wedge \delta \phi^A \right) \\ &= \int_M (\Upsilon_a \wedge \delta \theta^a + \Pi_a \wedge \delta T^a + \Phi_A \wedge \delta \phi^A).\end{aligned}$$

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- Matter part:

$$\delta S_m = \int_M (\Sigma_a \wedge \delta \theta^a + \Psi_A \wedge \delta \phi^A + \Omega_I \wedge \delta \chi^I).$$

General scalar-torsion gravity - field equations

- Local Lorentz invariance:

- Gravitational part:

$$\Upsilon^{[a} \wedge \theta^{b]} + \Pi^{[a} \wedge T^{b]} = 0 \quad \Leftrightarrow \quad \Delta^{[a} \wedge \theta^{b]} - \frac{1}{2} \overset{\bullet}{D} \Xi^{ab} = 0.$$

- Matter part: $\Sigma^{[a} \wedge \theta^{b]} = 0$.

General scalar-torsion gravity - field equations

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- Diffeomorphism invariance:

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- Field equations:

- Tetrad field equations:

$$\Delta_a + \Sigma_a = 0 \quad \Leftrightarrow \quad \Upsilon_a - \overset{\bullet}{D} \Pi_a + \Sigma_a = 0.$$

- Antisymmetric part \equiv connection field equations:

$$\overset{\bullet}{D} \Xi^{ab} = 0 \quad \Leftrightarrow \quad \overset{\bullet}{D} \Pi^{[a} \wedge \theta^{b]} + \Pi^{[a} \wedge T^{b]} = 0.$$

- Scalar field equations: $\Phi_A + \Psi_A = 0$.
 - Matter field equations: $\Omega_I = 0$.

$L(T, X, Y, \phi)$ theory - action

- Gravitational part of the action [MH, C. Pfeifer '18]:

$$S_g \left[\theta^a, \overset{\bullet}{\omega}{}^a{}_b, \phi^A \right] = \int_M L \left(T, X^{AB}, Y^A, \phi^A \right) \theta d^4x .$$

- Torsion scalar: $T = \frac{1}{2} T^\rho_{\mu\nu} S_\rho^{\mu\nu}$.
- Superpotential:

$$S_{\rho\mu\nu} = \frac{1}{2} (T_{\nu\mu\rho} + T_{\rho\mu\nu} - T_{\mu\nu\rho}) - g_{\rho\mu} T^\sigma_{\sigma\nu} + g_{\rho\nu} T^\sigma_{\sigma\mu} .$$

- Scalar field kinetic term: $X^{AB} = -\frac{1}{2} g^{\mu\nu} \phi^A_{,\mu} \phi^B_{,\nu}$.
- Kinetic coupling term: $Y^A = T_\mu^{\mu\nu} \phi^A_{,\nu}$.

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- Matter action variation expressed in components:

$$\delta S_m [\theta^a, \phi^A, \chi^I] = \int_M \left(\Theta_a{}^\mu \delta \theta^a{}_\mu + \vartheta_A \delta \phi^A + \varpi_I \delta \chi^I \right) \theta d^4x .$$

$L(T, X, Y, \phi)$ theory - field equations

- Symmetric part of tetrad equations:

$$\begin{aligned} \overset{\circ}{\nabla}_{(\mu} \left(L_{Y^A} \phi_{,\nu)}^A \right) - \overset{\circ}{\nabla}_\sigma \left(L_{Y^A} \phi_{,\rho}^A \right) g^{\rho\sigma} g_{\mu\nu} + L_{Y^A} \left(T_{(\mu\nu)}{}^\rho \phi_{,\rho}^A + T^\rho{}_{\rho(\mu} \phi_{,\nu)}^A \right) \\ - L g_{\mu\nu} - 2 \overset{\circ}{\nabla}_\rho \left(L_T S_{(\mu\nu)}{}^\rho \right) + L_T S_{(\mu}{}^{\rho\sigma} T_{\nu)\rho\sigma} - L_{X^{AB}} \phi_{,\mu}^A \phi_{,\nu}^B = \Theta_{\mu\nu}. \end{aligned}$$

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- Antisymmetric part of tetrad equations \equiv connection equations:

$$3\partial_{[\rho} L_T T^\rho{}_{\mu\nu]} + \partial_{[\mu} L_{Y^A} \phi_{,\nu]}^A - \frac{3}{2} L_{Y^A} T^\rho{}_{[\mu\nu} \phi_{,\rho]}^A = 0.$$

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- Scalar field equations:

$$g^{\mu\nu} \overset{\circ}{\nabla}_\mu \left(L_{Y^A} T^\rho{}_{\rho\nu} - L_{X^{AB}} \phi_{,\nu}^B \right) - L_{\phi^A} = \vartheta_A.$$

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- Matter field equations: $\varpi_I = 0$.

“Scalar-curvature”-like class - action

- Action [MH '18]:

- Gravitational part:

$$S_g \left[\theta^a, \overset{\bullet}{\omega}{}^a{}_b, \phi^A \right] = \frac{1}{2\kappa^2} \int_M \left[-\mathcal{A}(\phi)T + 2\mathcal{B}_{AB}(\phi)X^{AB} + 2\mathcal{C}_A(\phi)Y^A - 2\kappa^2\mathcal{V}(\phi) \right] \theta d^4x .$$

- Matter part:

$$S_m[\theta^a, \phi, \chi^I] = S_m^{\tilde{\jmath}} \left[e^{\alpha(\phi)} \theta^a, \chi^I \right] .$$

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- Free functions $\mathcal{A}, \mathcal{B}_{AB}, \mathcal{C}_A, \mathcal{V}, \alpha$ of scalar fields.

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- Free functions $\mathcal{A}, \mathcal{B}_{AB}, \mathcal{C}_A, \mathcal{V}, \alpha$ of scalar fields.
- $\mathcal{C}_A \equiv -\mathcal{A}_{,A} \Leftrightarrow$ theory reduces to scalar-curvature gravity.

“Scalar-curvature”-like class - field equation

- Symmetric part of the tetrad equations:

$$\begin{aligned} & (\mathcal{A}_{,A} + \mathcal{C}_A) S_{(\mu\nu)}{}^\rho \phi^A_{,\rho} + \mathcal{A} \left(\overset{\circ}{R}_{\mu\nu} - \frac{1}{2} \overset{\circ}{R} g_{\mu\nu} \right) + \left(\frac{1}{2} \mathcal{B}_{AB} - \mathcal{C}_{(A,B)} \right) \phi^A_{,\rho} \phi^B_{,\sigma} g^{\rho\sigma} g_{\mu\nu} \\ & - (\mathcal{B}_{AB} - \mathcal{C}_{(A,B)}) \phi^A_{,\mu} \phi^B_{,\nu} + \mathcal{C}_A \left(\overset{\circ}{\nabla}_\mu \overset{\circ}{\nabla}_\nu \phi^A - \overset{\circ}{\square} \phi^A g_{\mu\nu} \right) + \kappa^2 \mathcal{V} g_{\mu\nu} = \kappa^2 \Theta_{\mu\nu}, \end{aligned}$$

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- Scalar field equation:

$$\begin{aligned} & \frac{1}{2} \mathcal{A}_{,A} T - \mathcal{B}_{AB} \overset{\circ}{\square} \phi^B - \left(\mathcal{B}_{AB,C} - \frac{1}{2} \mathcal{B}_{BC,A} \right) g^{\mu\nu} \phi_{,\mu}^B \phi_{,\nu}^C \\ & + \mathcal{C}_A \overset{\circ}{\nabla}_\mu T_\nu{}^{\nu\mu} + 2\mathcal{C}_{[A,B]} T_\mu{}^{\mu\nu} \phi_{,\nu}^B + \kappa^2 \mathcal{V}_{,A} = \kappa^2 \alpha_{,A} \Theta. \end{aligned}$$

“Scalar-curvature”-like class - conformal transf.

- Conformal transformation and scalar field redefinition:

$$\bar{\theta}^a = e^{\gamma(\phi)} \theta^a, \quad \bar{e}_a = e^{-\gamma(\phi)} e_a, \quad \bar{\phi}^A = f^A(\phi).$$

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- Transformation of geometry:

$$\bar{T} = e^{-2\gamma} \left(T + 4\gamma_{,A} Y^A + 12\gamma_{,A}\gamma_{,B} X^{AB} \right), \quad \bar{\phi}^A = f^A,$$

$$\bar{X}^{AB} = e^{-2\gamma} \frac{\partial \bar{\phi}^A}{\partial \phi^C} \frac{\partial \bar{\phi}^B}{\partial \phi^D} X^{CD}, \quad \bar{Y}^A = e^{-2\gamma} \frac{\partial \bar{\phi}^A}{\partial \phi^B} \left(Y^B + 6\gamma_{,C} X^{BC} \right),$$

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- Transformation of parameter functions to preserve action:

$$\mathcal{A} = e^{2\gamma} \bar{\mathcal{A}},$$

$$\mathcal{B} = e^{2\gamma} \left(\bar{\mathcal{B}} f'^2 - 6\bar{\mathcal{A}} \gamma'^2 + 6\bar{\mathcal{C}} f' \gamma' \right),$$

$$\mathcal{C} = e^{2\gamma} (\bar{\mathcal{C}} f' - 2\bar{\mathcal{A}} \gamma'),$$

$$\mathcal{V} = e^{4\gamma} \bar{\mathcal{V}},$$

$$\alpha = \bar{\alpha} + \gamma.$$

“Scalar-curvature”-like class - invariants

- Quantities invariant under conformal transformations γ :

- “Scalar” quantities:

$$\mathcal{I}_1 = \frac{e^{2\alpha}}{\mathcal{A}}, \quad \mathcal{I}_2 = \frac{\mathcal{V}}{\mathcal{A}^2}.$$

- “Covector” quantities:

$$\mathcal{H}_A = \frac{\mathcal{C}_A + \mathcal{A}_{,A}}{2\mathcal{A}}, \quad \mathcal{K}_A = \frac{\mathcal{C}_A + 2\alpha_{,A}\mathcal{A}}{2e^{2\alpha}}.$$

- “Metric” quantities:

$$\begin{aligned}\mathcal{F}_{AB} &= \frac{2AB\mathcal{B}_{AB} - 6\mathcal{A}_{,(A}\mathcal{C}_{B)} - 3\mathcal{A}_{,A}\mathcal{A}_{,B}}{4\mathcal{A}^2}, \\ \mathcal{G}_{AB} &= \frac{\mathcal{B}_{AB} - 6\alpha_{,(A}\mathcal{C}_{B)} - 6\alpha_{,A}\alpha_{,B}\mathcal{A}}{2e^{2\alpha}}.\end{aligned}$$

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- Covariance under scalar field redefinitions:

$$\bar{\mathcal{I}}_{1,2} = \mathcal{I}_{1,2}, \quad (\bar{\mathcal{H}}, \bar{\mathcal{K}})_A = \frac{\partial \phi^B}{\partial \bar{\phi}^A} (\mathcal{H}, \mathcal{K})_B, \quad (\bar{\mathcal{F}}, \bar{\mathcal{G}})_{AB} = \frac{\partial \phi^C}{\partial \bar{\phi}^A} \frac{\partial \phi^D}{\partial \bar{\phi}^B} (\mathcal{F}, \mathcal{G})_{CD}.$$

“Scalar-curvature”-like class - special frames

- Jordan frame: minimal coupling to matter.

$$\mathcal{A}^{\mathfrak{J}} = \frac{1}{\mathcal{I}_1}, \quad \mathcal{B}_{AB}^{\mathfrak{J}} = 2\mathcal{G}_{AB}, \quad \mathcal{C}_A^{\mathfrak{J}} = 2\mathcal{K}_A, \quad \mathcal{V}^{\mathfrak{J}} = \frac{\mathcal{I}_2}{\mathcal{I}_1^2}, \quad \alpha^{\mathfrak{J}} = 0.$$

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$$\mathcal{A}^{\mathfrak{J}} = \frac{1}{\mathcal{I}_1}, \quad \mathcal{B}_{AB}^{\mathfrak{J}} = 2\mathcal{G}_{AB}, \quad \mathcal{C}_A^{\mathfrak{J}} = 2\mathcal{K}_A, \quad \mathcal{V}^{\mathfrak{J}} = \frac{\mathcal{I}_2}{\mathcal{I}_1^2}, \quad \alpha^{\mathfrak{J}} = 0.$$

- Einstein frame: no coupling to torsion scalar.

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- “Debraiding frame” (for $\mathcal{H}_{[A,B]} \equiv 0$): minimal coupling to torsion.

$$\left(\ln \mathcal{A}^{\mathfrak{D}} \right)_{,A} = 2\mathcal{H}_A, \quad \left(\ln \mathcal{B}^{\mathfrak{D}} \right)_{B,C}^A = [\ln (\mathcal{F} + 3\mathcal{H} \otimes \mathcal{H})]_{}^A_{B,C} + 2\delta_B^A \mathcal{H}_C,$$
$$\textcolor{red}{\mathcal{C}_A^{\mathfrak{D}} = 0}, \quad \left(\ln \mathcal{V}^{\mathfrak{D}} \right)_{,A} = (\ln \mathcal{I}_2)_A + 4\mathcal{H}_A, \quad \alpha^{\mathfrak{D}}_{,A} = \mathcal{I}_1 \mathcal{K}_A.$$

Scalar-torsion gravity without derivative coupling

- Gravitational action [MH, L. Järv, U. Ualikhanova '18]:

$$S = \frac{1}{2\kappa^2} \int_M \left[f(T, \phi) + Z_{AB}(\phi) g^{\mu\nu} \phi_{,\mu}^A \phi_{,\nu}^B \right] \theta d^4x + S_m[\theta^a, \chi^I].$$

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- Field equations:

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$$\partial_{[\rho} f_T T_{\mu\nu]}^\rho = 0 \quad \Leftrightarrow \quad \partial_\mu f_T \left[\partial_\nu (\theta e_{[a}{}^\mu e_{b]}{}^\nu) + 2\theta e_c{}^{[\mu} e_{[a}{}^{\nu]} \overset{\bullet}{\omega}_{b]\nu}^c \right] = 0.$$

- Scalar field equation:

$$f_{\phi^A} - (2Z_{AB,\phi^C} - Z_{BC,\phi^A}) g^{\mu\nu} \phi_{,\mu}^B \phi_{,\nu}^C - 2Z_{AB} \overset{\circ}{\square} \phi^B = 0.$$

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- Equation possesses generic solutions with spacetime symmetries.

Conclusion

- Summary:
 - Four new classes of scalar-torsion theories.

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$$L(T, X, Y, \phi)$$

no derivative coupling

“scalar-curvature-like”

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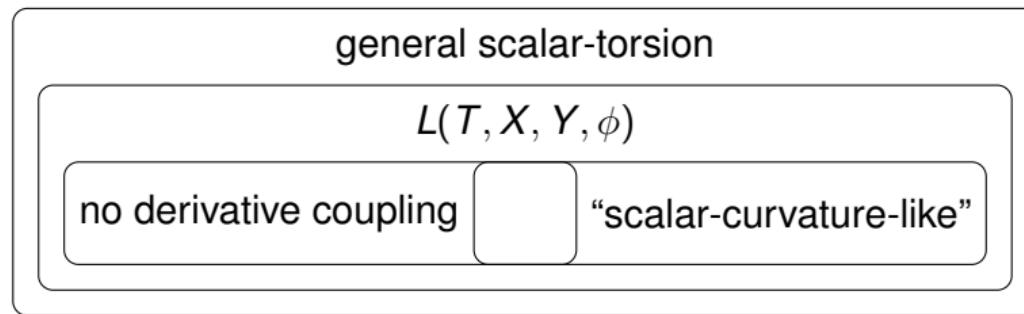
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- Outlook - analyze various aspects of these theories:

- Cosmological dynamics / dynamical systems analysis.
- Post-Newtonian limit.
- Gravitational waves - speed and polarisations.
- Hamiltonian formulation - degrees of freedom.

Literature

- MH, L. Järv, U. Ualikhanova; Covariant formulation of scalar-torsion gravity; Phys. Rev. D **97** (2018) 104011 [arXiv:1801.05786].
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